

# NONCOMMUTATIVE BATALIN-VILKOVISKY GEOMETRY AND MATRIX INTEGRALS

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ABSTRACT. We describe matrix model associated naturally with any solution to quantum master equation of noncommutative Batalin-Vilkovisky geometry. In the simplest example we get the Kontsevich model of 2-dimensional gravity.

## 1. INTRODUCTION.

We describe matrix model associated naturally with solution to quantum master equation of noncommutative Batalin-Vilkovisky geometry. The simplest example is the Kontsevich model of 2-dimensional gravity.

Notations. We work in the tensor category of super vector spaces, over an algebraically closed field  $k$ ,  $\text{char}(k) = 0$ . Let  $V = V^{1|0} \oplus V^{0|1}$  be a super vector space. We denote by  $\bar{\alpha}$  the parity of an element  $\alpha$  and by  $\Pi V$  the super vector space with inversed parity. Sometimes we may also assume that  $V$  has a differential, we leave to the interested reader to make the appropriate modifications, if they are not mentioned explicitly. For a finite group  $G$  acting on a vector space  $U$ , we denote via  $U^G$  the space of invariants with respect to the action of  $G$ .

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## 2. NONCOMMUTATIVE BATALIN-VILKOVISKY GEOMETRY.

2.1. **Even inner products.** Let  $B : V^{\otimes 2} \rightarrow k$  be an even symmetric inner product on  $V$ :

$$B(x \otimes y) = (-1)^{\bar{x}\bar{y}} B(y \otimes x)$$

The inner product  $B$  corresponds to an even symplectic structure on  $\Pi V$ . We introduced in [B1] the space  $F = \bigoplus_n F_n$

$$(2.1) \quad F_n = ((\Pi V)^{\otimes n} \otimes k[\mathbb{S}_n]')^{\mathbb{S}_n}$$

where  $k[\mathbb{S}_n]'$  denotes the super  $k$ -vector space with the basis indexed by elements  $(\sigma, \rho_\sigma)$ , where  $\sigma \in \mathbb{S}_n$  is a permutation with  $i_\sigma$  cycles  $\sigma_\alpha$  and  $\rho_\sigma = \sigma_1 \wedge \dots \wedge \sigma_{i_\sigma}$ ,  $\rho_\sigma \in \text{Det}(\text{Cycle}(\sigma))$ ,  $\text{Det}(\text{Cycle}(\sigma)) = \text{Symm}^{i_\sigma}(k^{0|i_\sigma})$ , is one of the generators of

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the one-dimensional determinant of the set of cycles of  $\sigma$ , i.e.  $\rho_\sigma$  is an order on the set of cycles defined up to even reordering, and  $(\sigma, -\rho_\sigma) = -(\sigma, \rho_\sigma)$ . The group  $\mathbb{S}_n$  acts on  $k[\mathbb{S}_n]'$  by conjugation. The space  $F$  carries the naturally defined operator

$$\Delta(a_1 \otimes \dots \otimes a_n \otimes (\sigma, \rho_\sigma)) = B(a_{n-1} \otimes a_n)(a_1 \otimes \dots \otimes a_{n-2} \otimes (\pi_{n-1,n}(s_{n-1,n}\sigma), \rho_\delta))$$

where  $s_{n-1,n}$  is the transposition interchanging  $n-1$  and  $n$ ,  $\pi_{n-1,n}$  is the operator of erasing of elements  $n-1$  and  $n$  from the cycles of the representation, and  $\rho_\delta$  is the natural ordering on the cycles of resulting permutation  $\pi_{n-1,n}(\sigma s_{n-1,n})$ , see section 9 from loc.cit. The space  $F$  can be identified with the symmetric power of the parity transposed space of cyclic invariant tensors on  $\Pi V$

$$F = \text{Symm}(\oplus_j \Pi(\Pi V^{\otimes j})^{\mathbb{Z}/j\mathbb{Z}})$$

If one chooses a basis  $\{a_i\}$  in  $\Pi V$ , in which  $B$  has the form  $B(a_i, a_j) = b_{ij}$  then operator  $\Delta$  can be written as

$$\Delta = \sum_{ij} \gamma b_{ij} \frac{\partial^2}{\partial a_i \partial a_j}$$

where  $\gamma$  denotes the "dissection-gluing" operator acting on cycles, which corresponds to the symmetrisation of the operator  $(\sigma, \rho_\sigma) \rightarrow (\pi_{n-1,n}(s_{n-1,n}\sigma), \rho_\delta)$  acting on pairs  $(\sigma, \rho_\sigma)$ . The operator  $\gamma$  is odd since it changes the total number of cycles by one.

The operator  $\Delta$  is of the second order with respect to the multiplication on  $F$  and by the general formula of Batalin-Vilkovisky geometry it defines the odd Lie bracket on  $F$ , which we described in loc.cit., section 5, and which can be written as

$$\{x, y\} = \sum_{ij} \gamma b_{ij} \frac{\partial x}{\partial a_i} \frac{\partial y}{\partial a_j}$$

Let us denote by  $\tilde{\mathbb{S}}$  the  $\mathbb{Z}/2\mathbb{Z}$ -graded version of the twisted modular  $\mathcal{K}$ -operad  $\tilde{\mathfrak{S}}\mathbb{S}[t]$ . By (a  $\mathbb{Z}/2\mathbb{Z}$ -graded version) of the theorem 1 from [B1] the solutions of quantum master equation in  $F$

$$d_V S + \hbar \Delta S + \frac{1}{2} \{S, S\} = 0,$$

$S = \sum_{2(g-1)+n>0} \hbar^g S_{g,n}$ , are in one-to one correspondence with the structure of algebra over the  $\mathbb{Z}/2\mathbb{Z}$ -graded modular operad  $\mathcal{F}_\mathcal{K} \tilde{\mathbb{S}}$  on the vector space  $V$  with symmetric inner product  $B$ . Its genus zero part corresponds to the structure of  $A_\infty$ -algebra with the even invariant inner product on  $V$ .

**2.2. Odd inner products.** Let  $V = V^{1|0} \oplus V^{0|1}$  be a super vector space and  $B : V^{\otimes 2} \rightarrow \Pi k$  be an odd symmetric inner product,  $B(x \otimes y) = (-1)^{\bar{x}+\bar{y}} B(y \otimes x)$ . By definition  $B$  is an odd symplectic structure on  $V$ . The space  $F$  in this situation has components

$$(2.2) \quad F_n = (V^{\otimes n} \otimes k[\mathbb{S}_n])^{\mathbb{S}_n}$$

where  $k[\mathbb{S}_n]$  is the  $k$ -vector space with the basis indexed by  $\sigma \in \mathbb{S}_n$  and  $\mathbb{S}_n$  acts on  $k[\mathbb{S}_n]$  by conjugation. The space  $F$  carries the naturally defined operator

$$\Delta(a_1 \otimes \dots \otimes a_n \otimes \sigma) = (-1)^{\sum_{i=1}^{n-2} \bar{a}_i} B(a_{n-1} \otimes a_n)(a_1 \otimes \dots \otimes a_{n-2} \otimes (\pi_{n-1,n}(s_{n-1,n}\sigma))$$

where  $s_{n-1,n}$  is the transposition interchanging  $n-1$  and  $n$ ,  $\pi_{n-1,n}$  is the operator of erasing of elements  $n-1$  and  $n$  from the cycles of the representation, see section

9 from loc.cit. The space  $F$  can be identified with the symmetric power of the space of cyclic invariant tensors on  $V$

$$F = \text{Symm}(\oplus_i (V^{\otimes i})^{\mathbb{Z}/i\mathbb{Z}})$$

If one chooses a basis  $\{a_i\}$  in  $V$ , in which  $B$  has the form  $B(a_i, a_j) = b_{ij}$  then operator  $\Delta$  can be written as

$$\Delta = \sum_{ij} \gamma^{b_{ij}} \frac{\partial^2}{\partial a_i \partial a_j}$$

where  $\gamma$  denotes the "dissection-gluing" operator acting on cycles, which corresponds to the symmetrisation of the operator  $\sigma \rightarrow \pi_{n-1, n}(s_{n-1, n}\sigma)$  acting on permutations. The operator  $\gamma$  is even in this case since in the definition of  $F$  the parity change on the space of cyclic invariant tensors is not involved.

The operator  $\Delta$  is of the second order with respect to the multiplication on  $F$  and it defines the odd Lie bracket on  $F$ , which can be written as

$$\{x, y\} = \sum_{ij} \gamma^{b_{ij}} \frac{\partial x}{\partial a_i} \frac{\partial y}{\partial a_j}$$

The odd analogue corresponds to the modular operad, which is natural to denote  $\mathbb{S}$ , the *untwisted* analogue of  $\mathbb{Z}/2\mathbb{Z}$ -graded version of  $\tilde{\mathfrak{S}}\mathbb{S}[t]$ . The spaces  $\mathbb{S}((n))$  are simply the group algebras  $k[\mathbb{S}_n]$  with the action of  $\mathbb{S}_n$  by conjugation. The composition maps are defined by operators  $\pi_{f, f'}(s_{f, f'}\sigma)$  as in section 9 of [B1]. The Feynman transform  $\mathcal{F}\mathbb{S}[t]$  is a  $\mathcal{K}$ -twisted modular operad. The solutions of quantum master equation in  $F$  are in one-to one correspondence with the structure of algebra over the modular operad  $\mathcal{F}\mathbb{S}[t]$  on the vector space  $V$  with *odd* symmetric inner product  $B$ . The genus zero part of  $\mathcal{F}\mathbb{S}[t]$  action on  $V$  corresponds to the structure of  $A_\infty$ -algebra with the odd invariant inner product on  $V$ . It is straightforward generalisation of the theorem 2 from loc.cit. that the complexes  $\mathcal{F}\mathbb{S}[t]((0, \gamma, \nu))$  compute the cohomology of the compactified moduli spaces with coefficients in the local system  $H^i(\overline{\mathcal{M}}'_{\gamma, \nu}/\mathbb{S}_\nu, \mathcal{L})$ ,  $\mathcal{L} = \text{Det}(P_\Sigma)$ , where  $P_\Sigma$  is the set of marked points on the Riemann surface  $\Sigma$ .

### 3. MATRIX INTEGRALS.

We showed in [B1] that solutions to quantum master equation corresponding to a twisted modular operad  $\mathcal{P}$  define natural cocycles in the graph complex which is the Feynmann transform of  $\mathcal{P}$ . In this section we show that for  $\tilde{\mathbb{S}}$  and  $\mathbb{S}$  these cocycles can be naturally interpreted as the terms of Feynman diagram expansions of matrix integrals of  $\exp(S/\hbar)$ .

**3.1. Odd inner product.** Let us assume that

$$S(a_i, \hbar) = \sum_{2(g-1)+n>0} \hbar^g S_{n, g}(a_i)$$

$S_{n, g} \in F_n$ , be a solution to quantum master equation in the space (2.2). Let us consider the super vector space

$$M = \text{Hom}(V, \text{End}_k(U))$$

where  $U$  is a super vector space  $\dim U < \infty$ . If  $V = k^{(d|f)}$ , then  $M = (End_k(U))^{\times d} \times (\Pi End_k(U))^{\times f}$ . The supertrace gives a natural even symmetric inner product on the space  $End_k(U)$  :

$$Tr(X \otimes Y) = \Sigma_{\alpha\beta} (-1)^{\bar{\alpha}} X_{\beta}^{\alpha} Y_{\alpha}^{\beta}$$

$X, Y \in End_k(U)$ , it is invariant with respect to the action of the group  $GL(U)$ . It gives an extension to  $M$  of the odd symmetric inner product on  $Hom(V, k)$  dual to  $B$ . In coordinates it is given by

$$B^{-1}(X, Y) = \sum_{i,j,\alpha,\beta} (B^{-1})^{ij} (-1)^{\bar{\alpha}} X_{i,\beta}^{\alpha} Y_{j,\alpha}^{\beta}$$

Let us extend  $S$  to a function  $S_U$  on  $M$  in the obvious way, so that instead of each cyclically symmetric tensor we put the supertrace of the product of the corresponding matrices from  $End_k(U)$

$$\sum a_{i_1} \otimes \dots \otimes a_{i_k} \rightarrow Tr(X_{i_1} \dots X_{i_k})$$

and the product of cyclic words goes to the product of traces.

We claim that the tensors which we associated with  $S(a_i, \hbar)$  via the statistical sum model construction to a graph  $G$ , equipped with the marking corresponding to basis elements from  $\mathbb{MS}^{dual}((G))$ , are obtained naturally from the Feynman diagrams corresponding to the asymptotic expansion of the following integral

$$\log \int \exp \frac{1}{\hbar} \left( B^{-1}(X, \Xi) - \frac{1}{2} B^{-1}(X, X) + S_U(X) \right) dX / (2\pi\hbar)^{\frac{N}{2}}$$

(compare with the formula (0.1) from [GK]). This is immediate from the standard rules of the Feynman diagrammatics, in particular the combinatorics of the terms in  $S_U(X)$  reproduce exactly the data associated with vertices in the complex  $\mathbb{MS}^{dual}((G))$ , i.e. symmetric product of cyclic permutations and an integer number.

**3.2. Even inner product.** Let now  $S(a_i, \hbar) = \sum_{2(g-1)+n>0} \hbar^g S_{n,g}(a_i)$  is a solution to quantum master equation in the Batalin-Vilkovisky algebra 2.1. In this case the space  $M$  is the super vector space

$$M = Hom_k(\Pi V, \Pi End_k(U) \oplus End_k(U))$$

where  $U$  is a purely even vector space. The analogue of trace defines on the space  $\Pi End_k(U) \oplus End_k(U)$  an *odd* symplectic structure, which after choosing a basis  $\{e_{\alpha}\}$  in  $U$  can be written as

$$Tr(X \otimes Y) = \sum_{\alpha\beta} X_{\beta}^{\alpha} Y_{\alpha}^{\beta}$$

$X \in \Pi End_k(U)$ ,  $Y \in End_k(U)$ . Notice that there are natural isomorphisms

$$\Pi End_k(U) = Hom_k(\Pi U, U) = Hom_k(U, \Pi U)$$

and also

$$Hom_k(\Pi U, \Pi U) = End_k(U)$$

Therefore elements of the space  $\Pi End_k(U) \oplus End_k(U)$  can be naturally composed so that this space is in fact an associative algebra with odd invariant inner product.

We have an extension to  $M$  of the even symmetric inner product on  $Hom(V, k)$  dual to  $B$ :

$$B^{-1}(X, Y) = \sum_{i,j,\alpha,\beta} (B^{-1})^{ij} (-1)^{\bar{\alpha}} X_{i,\beta}^{\alpha} Y_{j,\alpha}^{\beta}$$

$X, Y \in \Pi\text{End}_k(U) \oplus \text{End}_k(U)$ .

We extend as above the solution  $S$  to a function  $S_U$  on  $M$  in the obvious way, so that instead of each cyclically symmetric tensor in  $(\Pi V)^{\otimes k}$  we put the trace of the product of the corresponding elements from  $(\Pi\text{End}_k(U) \oplus \text{End}_k(U))$

$$\sum a_{i_1} \otimes \dots \otimes a_{i_k} \rightarrow \text{Tr}(X_{i_1} \cdot \dots \cdot X_{i_k})$$

and the product of cyclic words goes to the product of traces.

Then the standard rules of the Feynman diagrammatics imply immediately that the asymptotic expansion of the integral

$$\log \int \exp \frac{1}{\hbar} \left( B^{-1}(X, \Xi) - \frac{1}{2} B^{-1}(X, X) + S_U(X) \right) dX / (2\pi\hbar)^{\frac{N}{2}}$$

$X, \Xi \in M$ , give the statistical sum tensors on graphs which were associated with the solution  $S$  in [B1].

**Example 1.** *Let us consider the case of even one-dimensional space  $V$ . It is easy to see that in this case any linear combination of cyclic words of  $X^3, X^5, \dots, X^{2n+1}$  is a solution of quantum master equation. The corresponding matrix integral reproduces the matrix integrals of [K].*

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